

Dedicated to Professor Bernhard Wunderlich on the occasion of his 65th birthday

## INSTANTANEOUS STRAIN RECOVERY ELASTICITY AND ENDOTHERMIC HEAT CHANGE IN NYLON-6 MONOFILAMENTS AND SILICONE RUBBERS

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### Abstract

The instantaneous elastic moduli for a nylon-6 monofilament were derived on strain recoveries right after creep, stress relaxation, and rapid elongation,  $E_c$ ,  $E_s$  and  $E_e$ , respectively. It was found that during strain recovery  $E_s (> E_c)$  and  $E_e$  increase monotonically with increasing load,  $m_1$ , on the sample. The extrapolated value of  $E_s$  at  $m_1 = 0$  g is almost equal to Young's modulus, 4.06 GPa. The value of  $E_e$  also increased with increasing  $m_1$ , and at  $m_1 = 600$  g ( $1.93 \text{ t cm}^{-2}$ ) reached about 14 GPa. The endothermic heat change right after creep, stress relaxation or rapid elongation,  $Q$ , was negligibly small. For comparison,  $E_s$ ,  $E_e$  and  $Q$  were also investigated for silicone rubber. It was found that  $E_s$  (53.8 MPa at the draw ratio  $D = 1.2$ ) decreased abruptly at  $D = 1.3$ . In the range of  $D = 1.4$ – $1.9$ ,  $E_s$  was only 22.6 MPa. In the case of stress relaxation,  $Q$  increased with increasing  $D$  from  $4 \text{ J mol}^{-1}$  (at  $D = 1.2$ ) to  $56 \text{ J mol}^{-1}$  (at  $D = 1.9$ ). Furthermore  $E_e$  (5.58 MPa at  $m_1 = 133.8$  g ( $429.4 \text{ kg cm}^{-2}$ )) increased gradually with increasing  $m_1$  and attained 16.6 MPa at  $m_1 = 548.4$  g ( $1.76 \text{ t cm}^{-2}$ ). In the case of creep,  $Q$  was in the range of 0– $11.5 \text{ J mol}^{-1}$  and larger when larger loads,  $m_2$ , were removed during the later stages of creep.

**Keywords:** instantaneous strain recovery elasticity, nylon-6 monofilament, silicone rubber

### Introduction

For highly oriented monofilaments [1, 2] of nylon-6 and polyvinylidenechloride and polypropylene films [3], the energetics of the strain recovery [4] has been observed by rapid removal of part or all of the load on a sample during creeping. The initial velocity of shrinkage was linked to the elastic modulus,  $E$ , of the instantaneous strain recovery [1–3]. In this study the behavior of instantaneous strain recovery of nylon-6 monofilaments right after creep, stress relaxation, or rapid elongation was observed. From the initial velocity of shrinkage, the elastic moduli of instantaneous strain recovery for the cases of creep, stress relaxation and rapid elongation were derived,  $E_c$ ,  $E_s$  and  $E_e$ , respectively. The endothermic heat,  $Q$ , right after rapid removal of a load on a

sample was negligible, indicating energy elasticity. For comparison,  $Q$ ,  $E_s$  and  $E_c$  were investigated for silicone rubber with entropy elasticity [5]. The rubber is exothermic when stretched rapidly and endothermic when the stress is released [6].

## Theoretical treatment

### *Elastic modulus of instantaneous strain recovery*

The behavior of instantaneous strain recovery is observed by rapid removal of part or all of the load on a sample during creep, stress relaxation, or rapid elongation. According to Newton's law, the kinetic equation of the initial behavior is given by:

$$m_1 x'' = -C(x_0 - x) + m_1 G - T \left( \frac{dS}{dx} \right)_{PT} \quad (1)$$

where  $x$  and  $x_0$  are the shrinkage and the final shrinkage after the rapid removal of a load from the sample, respectively;  $m_1$  is the load left on the sample;  $C$  is Hooke's constant;  $G$  is the gravitational acceleration;  $TdS/dx$  is a force due to the entropy elasticity [5] (here  $dx = -dL$ ,  $L$  is the sample length,  $S$  is the entropy);  $-TdS$  is related to the endothermic heat change caused by the instantaneous strain recovery. The third term on the right hand side of Eq. (1), which was negligibly small for a highly oriented nylon-6 monofilament, is predominant for rubbers and should be negligible for ideal energy elasticity. The solution,  $x$ , of Eq. (1) without the third term is given by [1, 3]:

$$x = \alpha_1 \exp \left\{ \left( \frac{C}{m_1} \right)^{1/2} t \right\} + \alpha_2 \exp \left\{ - \left( \frac{C}{m_1} \right)^{1/2} t \right\} + \frac{m_2 G}{C} \quad (2)$$

According to the conditions of  $x=0$  at  $t=0$  and  $x=m_2 G/C$  at  $t=\infty$ ,  $\alpha_1$  and  $\alpha_2$  should be 0 and  $-m_2 G/C$ , respectively; where  $t$  is the time and  $m_2$  is the load on a sample removed during creep, stress relaxation, or rapid elongation. Accordingly, Eq. (2) is rewritten as:

$$x = \left( \frac{m_2 G}{C} \right) \{ 1 - \exp(-qt) \}, \quad q = \left( \frac{C}{m_1} \right)^{1/2} \quad (3)$$

From Eq. (3),  $E$  is derived [1, 3]:

$$E = \left( \frac{CL}{\sigma} \right) = \left( \frac{L}{\sigma} \right) \left( \frac{1}{m_1} \right) \left( \frac{Gm_2}{V} \right)^2 \quad (4)$$

where  $V$  is the initial velocity of strain recovery and  $\sigma$  is the cross section of the sample. In the case of silicone rubber, the shrinkage velocity increased gradually and then attained a top velocity reflecting the entropy elasticity. If it is hypothesized that the top velocity is attributed to only energy elasticity, the top velocity corresponds to a simple harmonic motion, and  $C$  can be derived from Eq. (4) by choosing  $V$  at the top velocity. One can then calculate  $Q$  by:

$$Q = \frac{\int_0^{t_1} m_1 x'' x' dt + (C x_0 - m_1 G) x_1 - \frac{1}{2} C x_1^2}{M} \quad (5)$$

where  $t_1$  is the time when the top velocity is attained,  $x_1$  is  $x$  at  $t = t_1$  and  $M$  is the molecular weight of the sample.

## Experimental

### Sample

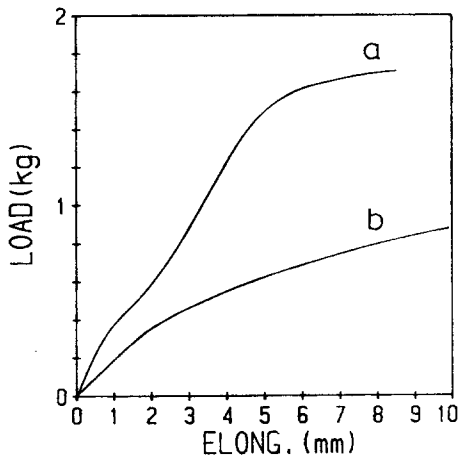
A nylon-6 monofilament (320 D) with the density  $1.141 \text{ g cm}^{-3}$  at  $25^\circ\text{C}$ , a crystallinity 30.0% and a degree of crystalline orientation of 0.94 was used. The second sample was a commercial silicone rubber in form of strips (thickness: 0.1 cm and 0.05 cm) with the width 0.3 cm.

### Young's modulus

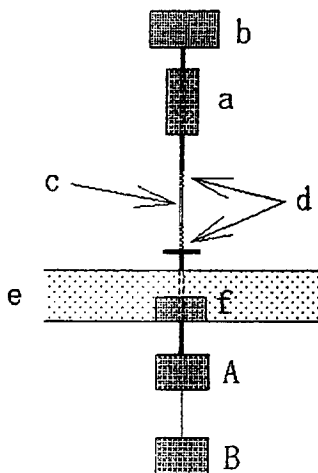
The measurement of load-elongation curves was carried out at room temperature of about  $20^\circ\text{C}$  and 35–45%RH. Young's moduli of 4.06 GPa for a nylon-6 monofilament (sample length: 2.5 cm), 2.81 MPa for a silicone rubber (sample length: 1.0 cm and thickness: 0.1 cm), and 3.17 MPa for another silicone rubber (sample length: 2.0 cm and thickness: 0.05 cm) were obtained from the initial slope of a load-elongation curve. The elongation rate was  $10 \text{ cm min}^{-1}$ . The dependence of Young's modulus on the elongation rate was not reproducible. Figure 1 shows the representative load-elongation curves for the nylon-6 monofilament and the silicone rubber.

### Instantaneous strain recovery right after rapid elongation

Figures 2 and 3 show the apparatus and a representative behavior at  $m_1 = 18.5 \text{ g}$  for a nylon-6 monofilament. In this measurement there is no load B in Fig. 2. All measurements of strain recovery were carried out at atmospheric pressure, about  $20^\circ\text{C}$  and 35–45%RH. The procedure of measurement was as follows: First a weight (a) is attached to the electromagnet (b) and the sample (c) is held by the grips (d) under tension of weight A. Then the weight (a) is dropped by switching off the electromagnetic. After the weight (a) reaches the



**Fig. 1** The load-elongation curves for a nylon-6 monofilament (a) (sample length: 2.5 cm) and a silicone rubber (b) (sample length: 1.0 cm, thickness: 0.1 cm). Elongation rate:  $10 \text{ cm min}^{-1}$



**Fig. 2** Apparatus for measuring the strain recovery behavior. a: weight to be dropped, b: electromagnetic, c: sample, d: grips, e: iron plate, f: displacement meter, A and B: loads

lower grip (d), the combination of the weight (a) and the lower grip (d) collides with the iron plate (e) and stretches the sample. The weight (a) rebounds upward after pressing the lower grip (d) against the iron plate (e) and the load forced on the sample is removed, except for the prior load A contained in the weight of the lower grip parts. At this time the shrinkage of the sample towards the initial extension is started. The initial shrinkage behaviors detected by the displacement meter (f) are drawn on the display by a computer and  $E_c$  beyond point  $P$  is given by:

$$E_e = \left( \frac{CL}{\sigma} \right) = \left( \frac{Lm_1}{\sigma} \right) \left( \frac{V}{x_0} \right)^2 \quad (6)$$

with  $x_0 = Gm_2/C$ .

The initial elongation velocity of a nylon-6 monofilament by the weight (a) was at an average of  $106.7 \text{ cm s}^{-1}$ .

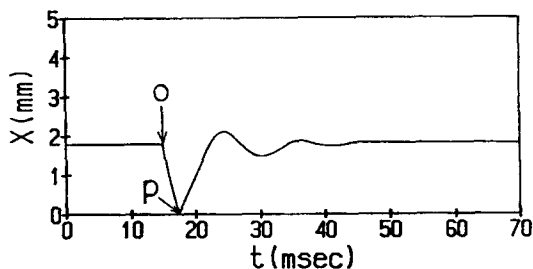


Fig. 3 Strain recovery behavior after rapid elongation for a nylon-6 monofilament.  $m_1 = 18.5 \text{ g}$ .  $X$ : shrinkage and  $t$ : time.  $O$  and  $P$ : beginning points of elongation and shrinkage

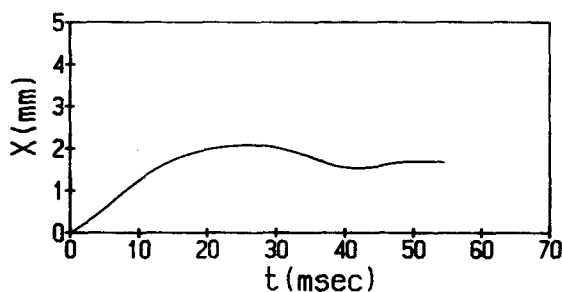


Fig. 4 Strain recovery behavior after creep for a nylon-6 monofilament.  $m_1 = 45.5 \text{ g}$  and  $m_2 = 498.6 \text{ g}$ .  $X$ : shrinkage and  $t$ : time

### *Instantaneous strain recovery right after creep*

The apparatus in Fig. 2 was used for nylon-6 monofilament and silicone rubber. First two loads A ( $=m_1$ ) and B ( $=m_2$ ) are suspended on a sample (c). After about 20 min, B was removed by cutting a thread connecting A and B. The initial behavior of a sample was detected by the displacement meter (f) and then was drawn on the display by a computer or recorded by a photocorder. Figures 4 and 5 show the initial behavior right after the rapid removal of B for a nylon-6 monofilament and a silicone rubber (thickness: 0.05 cm). As shown in Figs 4 and 5, for a nylon-6 monofilament an abrupt initial shrinkage was observed right after the rapid removal of  $m_2$ , and for a silicone rubber the shrinkage velocity increased gradually and then attained a top velocity.

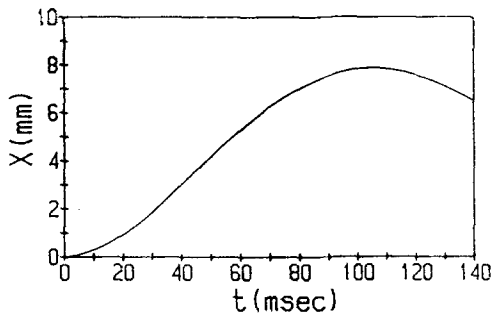


Fig. 5 Strain recovery behaviour after creep for a silicone rubber (thickness: 0.05 cm).  $m_1 = 133.8$  g and  $m_2 = 88.9$  g.  $X$ : shrinkage and  $t$ : time

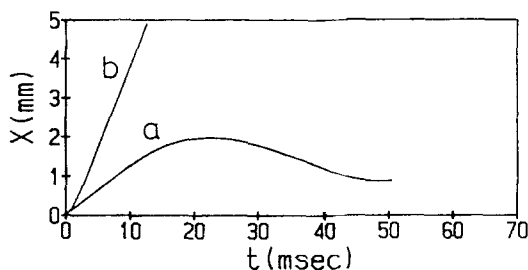


Fig. 6 Strain recovery behaviour after stress relaxation for a nylon-6 monofilament. a:  $m_1 = 135.9$  g and  $m_2 = 700$  g, b:  $m_1 = 45.5$  g and  $m_2 = 900$  g.  $X$ : shrinkage and  $t$ : time

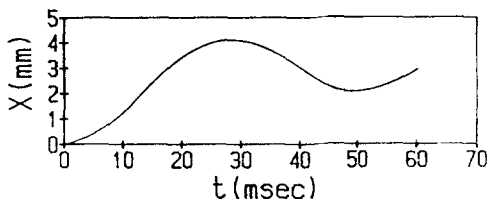


Fig. 7 Strain recovery behaviour after stress relaxation for a silicone rubber (thickness: 0.1 cm) at  $D = 1.40$ .  $m_1 = 45.5$  g and  $m_2 = 175$  g.  $X$ : shrinkage and  $t$ : time

### *Instantaneous strain recovery right after stress relaxation*

In this measurement, a load cell was installed instead of an electromagnet (b). First by two loads A ( $=m_1$ ) and B, a sample (c) is elongated to a constant length; the lower grip (d) is attached to the iron plate (e) by the load B of a heavy weight. Next when the relaxation load arrived at an appointed load, the thread connecting A and B is cut with scissors. The initial behavior under a load A right after the removal of B is drawn on the display by a computer and  $E_s$  was calculated from Eq. (4), where  $m_2$  is the relaxation load detected by the load cell just before the removal of B. Figures 6 and 7 show the initial behaviors right after the rapid removal of B during the stress relaxation for a nylon-6 monofila-

ment and a silicone rubber (thickness: 0.1 cm), respectively. The initial behaviors in these samples are similar to the cases of instantaneous strain recovery after creep.

## Results and discussion

### Nylon-6 monofilament

#### $E_e$ and $E_s$

Figure 8 shows the relationships between  $E$  ( $E_e$ ,  $E_s$  and  $E_c$  (described in the next section)) and  $m_1$  for a nylon-6 monofilament. As shown in Fig. 8,  $E_e$  and  $E_s$  ( $> E_e$ ) increase with increasing  $m_1$ , reflecting the change of the chain cohesive state. The dependence of  $E_s$  on  $m_2$  ( $=700$  g,  $800$  g and  $900$  g) was not found to be reproducible. The extrapolated value of  $E_s$  at  $m_1=0$  g was almost equal to Young's modulus ( $4.06$  GPa). The result of  $E_s > E_e$  may be attributed to the difference of relaxation times related to the relative cohesive state.

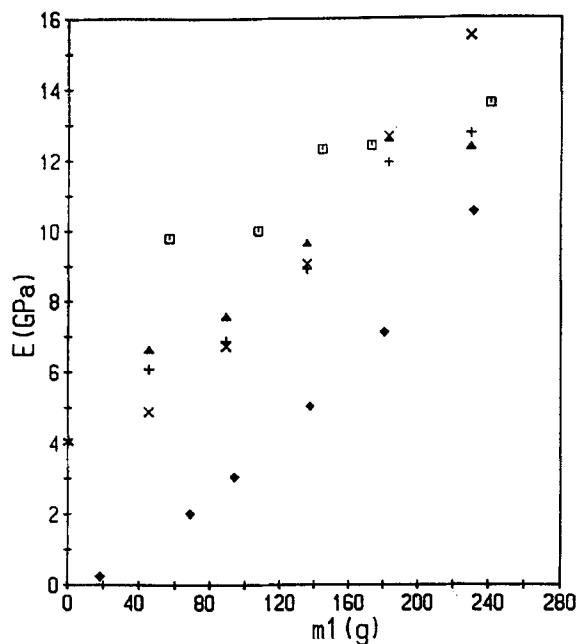


Fig. 8 Relationship between  $E$  ( $E_e$ ,  $E_s$  and  $E_c$ ) and  $m_1$  for a nylon-6 monofilament.  $E_e$ : ◆;  $E_s$ : × ( $m_2=700$  g), + ( $m_2=800$  g), ▲ ( $m_2=900$  g);  $E_c$ : □; Young's modulus: \*

#### $E_c$

Figure 9 shows the relationship between  $V$  and  $m_2$  in the case of creep for a nylon-6 monofilament. When  $m_1$  is  $56.9$  g,  $107.7$  g or  $240.5$  g, the values of  $V$

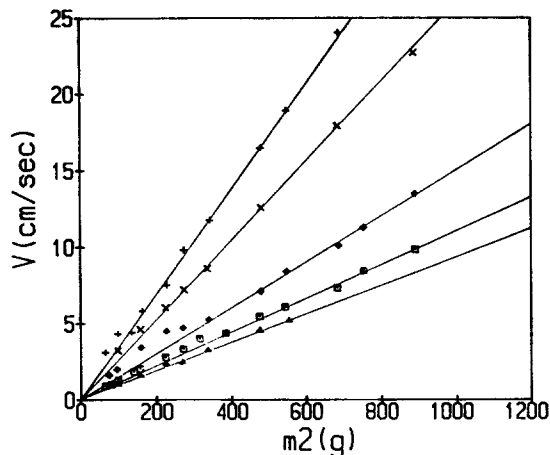


Fig. 9 Relationship between  $V$  and  $m_2$  for a nylon-6 monofilament.  $m_1$  (+: 56.9 g, x: 107.7 g, ♦: 240.5 g, ■: 423.5 g, ▲: 605.2 g)

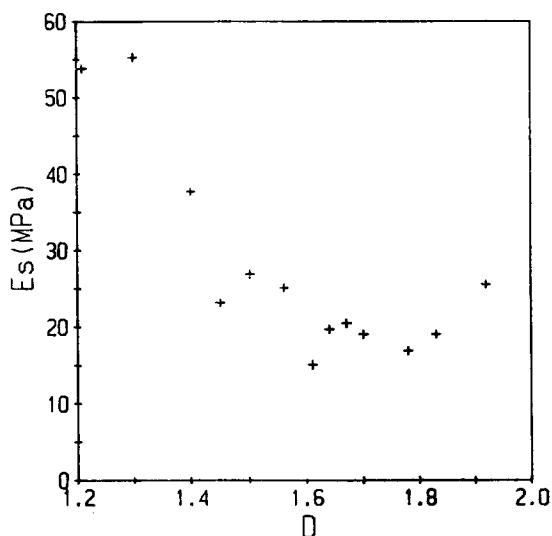


Fig. 10 Relationship between  $E_s$  and  $D$  for a silicone rubber (thickness: 0.1 cm)

in the lower range of  $m_2$  are a little larger than the straight solid line of extrapolation. The sum of  $m_1$  and the minimum value of  $m_2$  in the experimental straight line was in the yield load range of 200–500 g on the load – elongation curve. From the slope of each straight line,  $E_c$  was calculated using Eq. (4). The relationship between  $E_c$  and  $m_1$  is shown in Fig. 8.  $E_c$  increased with increasing  $m_1$  and reached about 14 GPa at  $m_1=600$  g (omitted in Fig. 8). Also in this case it is expected that a dense cohesive state could be introduced with increasing  $m_1$ . The value of  $E_c$  at  $m_1=173$  g is close to  $E_s$ .



## Silicone rubber

### $E_s$ and $Q$

Figure 10 shows the relationship between  $E_s$  and the draw ratio,  $D$ , for a silicone rubber (thickness: 0.1 cm). The value of  $E_s$  (53.8 MPa at  $D=1.2$ ) decreased abruptly over  $D=1.3$ . In the range of  $D = 1.4-1.9$ ,  $E_s$  was an average of 22.6 MPa. Figure 11 shows the relationship between  $Q$  and  $D$  for a silicone rubber (thickness: 0.1 cm).  $Q$  was 4–56 J mol<sup>-1</sup> in the range of  $D=1.2-1.9$ , and larger as  $D$  increased.

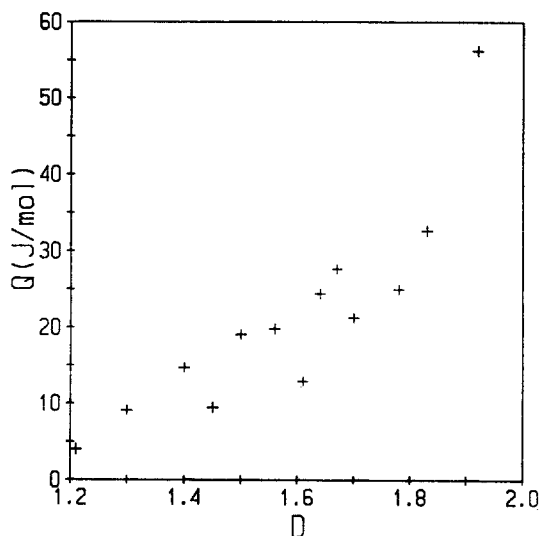
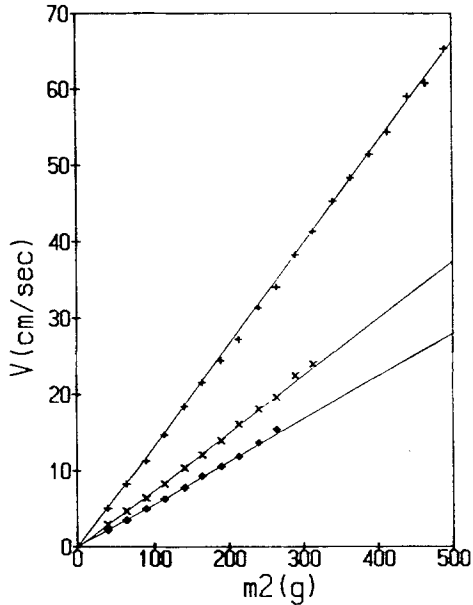


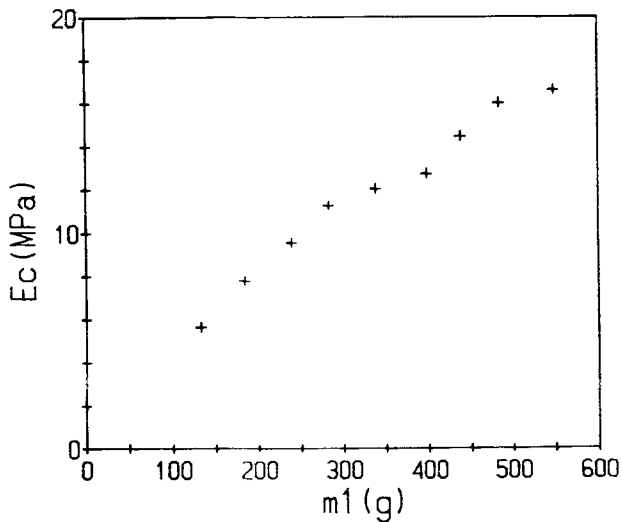
Fig. 11 Relationship between  $Q$  and  $D$  for a silicone rubber (thickness: 0.1 cm)

### $E_c$ and $Q$

Figure 12 shows the relationship between  $V$  and  $m_2$  in the case of creep for a silicone rubber (thickness: 0.05 cm). For each  $m_1$  (133.8 g, 238.2 g and 337.1 g) a straight line was obtained. From the slope of each straight line,  $E_c$  was calculated using Eq. (4). Figure 13 shows the relationship between  $E_c$  and  $m_1$ . The modulus  $E_c$  increased with increasing  $m_1$ , corresponding to the irreversible change of the cohesive state; the decrease of the recovery ratio to the initial length by increasing  $m_1+m_2$ . The recovery ratio was 80.0% at  $m_1+m_2=600$  g, 90.7% at  $m_1+m_2=450$  g, 96.3% at  $m_1+m_2=300$  g and 99.3% at  $m_1+m_2=150$  g. Figure 14 shows the relationship between  $Q$  and  $m_2$  at each  $m_1$  (238.2 g and 337.1 g). The value of  $Q$  was in the range of 0–11.5 J mol<sup>-1</sup> and larger for larger  $m_2$ . Here let's hypothesize that the highly stretched state of a silicone rubber is equivalent to its frozen (glassy) state in the entropy change [5]. The heat capacity difference  $\Delta C_p$  between a liquid and a glass at the glass

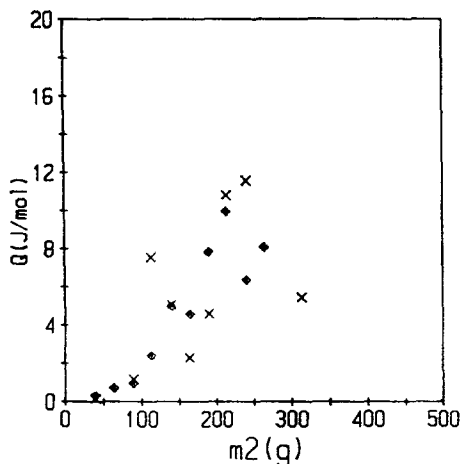


**Fig. 12** Relationship between  $V$  and  $m_2$  for a silicone rubber (thickness: 0.05 cm)  
 $m_1$  (+: 133.8 g,  $\times$ : 238.2 g and  $\blacklozenge$ : 337.1 g)



**Fig. 13** Relationship between  $E_c$  and  $m_1$  for a silicone rubber (thickness: 0.05 cm)

transition temperature,  $T_g$ , is  $27.7 \text{ J (K mol)}^{-1}$  for polydimethyl-siloxane ( $T_g = 146 \text{ K}$ ) [7]. Therefore, when the stretched state is released, a rise in temperature,  $\Delta T = 0.42 \text{ K}$  ( $= Q/\Delta C_p$ ,  $Q = 11.5 \text{ J mol}^{-1}$ ) at  $m_1 = 238.2 \text{ g}$  and



**Fig. 14** Relationship between  $Q$  and  $m_2$  at each  $m_1$  (238.2 g (x), and 337.1 g (◆)) for a silicone rubber (thickness: 0.05 cm)

$m_2=239.9$  g in the case of creep (Fig. 14) and  $\Delta T=2.03$  K ( $=Q/\Delta C_p$ ,  $Q=56.3$  J mol $^{-1}$ ) in the case of stress relaxation at  $D=1.92$  (Fig. 11), is expected. These values of  $\Delta T$  are reasonable comparing with 2.06 K for Neoprene at  $D=2.0$  and 0.42 K for Latex at  $D=2.0$  [6].

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